



S velocity reversal in the mantle transition zone

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[1] The transition zone differs from the rest of the mantle by high positive P and S wave velocity gradients. However, by applying S receiver function technique to recordings of about 50 globally distributed stations, in 7 regions we obtain evidence of negative discontinuity at a depth around 500 km. Four of these regions correspond to the well known hotspots. No other hotspots are sampled by our observations at this depth. Our modeling suggests that the negative discontinuity is hardly an effect of anisotropy of wadsleyite. More likely, it is caused by isotropic S velocity reduction of around 0.2 km/s. Origin of the low velocity can be related to enhanced water content in wadsleyite. **Citation:** Vinnik, L., and V. Farra (2006), S velocity reversal in the mantle transition zone, *Geophys. Res. Lett.*, 33, L18316, doi:10.1029/2006GL027120.

1. Introduction

[2] The mantle transition zone (MTZ) differs from the rest of the mantle by high positive P and S wave velocity gradients, which are caused mainly by a series of phase transformations. This depth range is notable by a low resolution and poor correlation of the global tomographic models. To understand processes in the MTZ, seismic studies with a better resolution are required. Application of S receiver function (SRF) technique [Farra and Vinnik, 2000] provided first evidence of a negative discontinuity in the upper MTZ in the vicinity of two hotspots: Afar [Vinnik *et al.*, 2004] and Iceland [Vinnik *et al.*, 2005]. Here we present a new series of seismic observations related to this discontinuity.

[3] Observations of the negative discontinuity are important, because the low S wave velocity in the MTZ can be related to the presence of water, which may have a large effect on processes in the mantle. Recent studies of the nominally anhydrous wadsleyite and ringwoodite indicate that they may contain up to about 2.0 wt% of H₂O [Hirschmann *et al.*, 2005]. The average content of water in the MTZ, though generally unknown, is, very likely, much lower than the maximum, but if it locally reaches 1 wt%, the S velocity of the hydrous phases becomes several percent lower than the anhydrous velocity [Smyth and Jacobsen, 2006].

[4] Like in the previous studies, our analysis is based on observations of the S_{dp} (converted from S to P at a depth *d*) phases that are detected using SRFs. The details of the technique can be found elsewhere [Farra and Vinnik, 2000], and here we present only the essentials. The three-component seismogram is decomposed into P and SV components. The SV axis corresponds to the principal S particle motion direction in the wave propagation plane. The P axis,

which is normal to SV in the same plane, is optimal for detecting the S_p converted phases, because the P component of the record is free from the S wave energy. The P components of many seismic recordings are deconvolved by the SV components of the S waves and stacked to reduce noise. Slant stacking facilitates detection of seismic phases, the slowness of which differs from that of the parent phase.

[5] The SRF method is relatively new, and its advantages and limitations still are debated. In the data there is a simple and robust indicator of quality of the SRF: a detection of the S_p phase from the global 410 km discontinuity implies that the method is sensitive enough to detect the S velocity contrast of around 0.2 km/s at this depth. Moreover, the largest depth sampled by the SRF depends on epicentral distance. The S_p phase from a negative discontinuity in the MTZ becomes observable at a smaller epicentral distance than the S_p phase from a positive discontinuity at this depth, and stacking in a broad distance range may result in a larger signal from the negative discontinuity.

[6] P receiver function (PRF) technique is an alternative method for mapping seismic boundaries in the MTZ. However, in the PRFs indications of a negative discontinuity were never found. Most of these studies were conducted at the periods from 2 to 5 seconds, ~ 3 times shorter than the periods of S receiver functions. A gradational boundary can be transparent at short periods in the PRF but detectable at longer periods in the SRF. Anelastic attenuation is another reason: the P_s phases can be a few times smaller than the S_p phases, owing to a high attenuation of the S waves relative to P. The most obvious advantage of the SRF over PRF stems from the fact that, unlike the P_s phases, the S_p phases from deep discontinuities are separated in time from reverberations of shallow origin.

2. Data and Results

[7] The data are extracted from a set of SRFs of more than 50 globally distributed seismograph stations (Figure 1) and portable groups of stations. Those stations, where we could not obtain SRFs of good quality, are not shown in Figure 1. The number of stacked receiver functions for each station in Figure 1 (at least several tens) is sufficient for obtaining a high signal/noise ratio: the normalized amplitudes of the signals of interest are around 0.03, whereas the RMS amplitude of noise is not more than 0.007 (Table 1). New observations of seismic signals from the negative discontinuity in the MTZ are made at stations BGCA and BNG in central Africa, TUC and CMB in the western USA, BJT in China, YKW3 and SCHQ in Canada and ALE and DSS in Afar.

[8] The recordings of neighboring stations BGCA and BNG were processed as a single data set. For our sign convention, negative polarity corresponds to positive discontinuity. The stack for the back azimuths around 78° (Figure 2a) is characteristic of the “normal” mantle: the only robust

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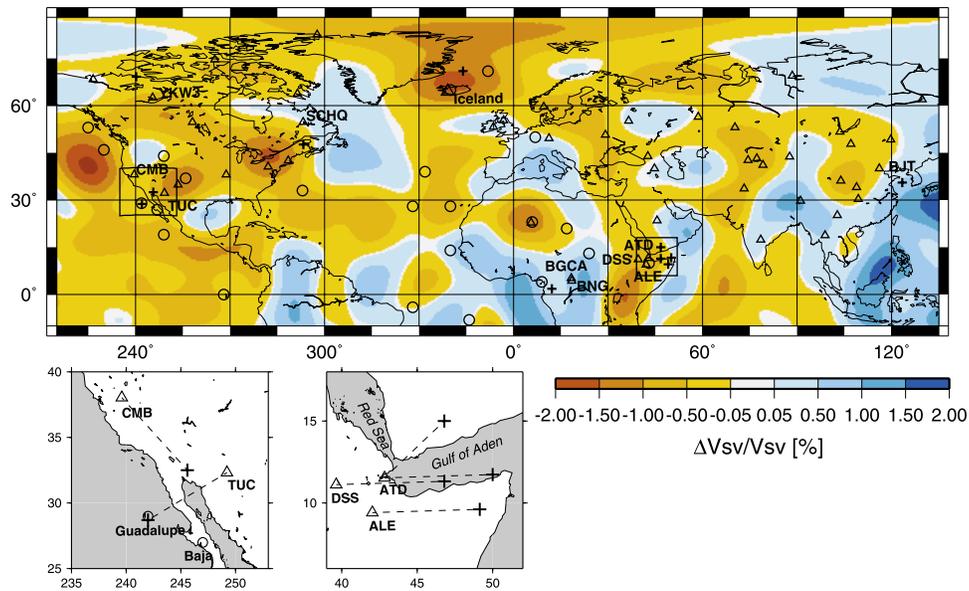


Figure 1. Map of the northern hemisphere with seismograph stations (triangles), piercing points (crosses), hotspots (circles, modified from *Duncan and Richards [1991]*) and S velocity at a depth of 525 km [*Ritsema and van Heist, 2000*]. Stations with indications of the negative discontinuity in the MTZ are marked by station codes. In Iceland the data are obtained at a group of stations [*Vinnik et al., 2005*]. At the bottom of the figure, Afar and the west of America are shown enlarged; seismograph stations are connected with the piercing points by dash lines. Outside the region shown on the map we have only a few seismograph stations (LBTB and BOSA in south Africa), NWA0 (Australia), SPA, SBA and DRV (Antarctica) and MA2 in easternmost Eurasia.

seismic phase in the time window of interest is S410p with negative polarity at -48 s. However, in the azimuths around 240° (Figure 2b), the dominant arrival has positive polarity at a time around -57 s. In our interpretation, this is the Sp phase from a negative discontinuity in the upper MTZ. The SRFs of stations CMB and TUC in differing azimuths (Figures 2c and 2d) sample neighboring regions in the MTZ (Figure 1), and they contain similar signals with positive polarity. To attain a high signal/noise ratio, the recordings of station ALE and DSS were processed jointly. The resulting stack (Figure 2e) is similar to that previously obtained in the same back azimuths (between 67° and 103°) at the neighboring station ATD [*Vinnik et al., 2004*]. Similar signals are present in the SRFs of stations BJT, YKW3 and SCHQ.

[9] Almost all SRFs in Figure 2 (except 2b) contain a clear S410p phase. The differences in the time of this phase reflect mainly differences in the average epicentral distance (Table 1). In a spherically symmetrical Earth, the theoretical differential slowness of the Sp phase from the upper MTZ are around 0.4 s/ $^\circ$. In most cases, the maximum amplitudes of the S410P

and of the Sp phase from the negative discontinuity occur near this slowness. The deviations are caused by noise and lateral heterogeneity of the Earth. In particular, a deviation of 0.2 s/ $^\circ$ can be caused by a tilt of around 1° of the converting interface. The SRFs in Figure 2 contain other interesting arrivals, but they are beyond the scope of our paper.

[10] To construct velocity models of the MTZ, we calculated synthetic SRFs [*Vinnik et al., 2004*], and searched for the preferred S velocity profiles by trial and error. The details of the models clearly cannot be determined uniquely, but at this stage we are interested only in the approximate depth range and magnitude of the S velocity reduction. We assume, that the 410-km discontinuity is located at its standard depth, and perturb IASP91 standard model [*Kennett and Engdahl, 1991*] above the 410 km discontinuity to obtain the model with the arrival time of S410p close to that in the actual data. Then we perturb S velocity in the MTZ to reproduce the seismic phase with positive polarity (see example in Figure 3a).

[11] We model the negative discontinuity by a gradational boundary 30 km thick, although our data cannot distinguish it

Table 1. Parameters of the Recordings

Station	Latitude, deg	Longitude, deg	Number of Events	Distance, deg	Azimuth, deg	Noise	Depth, km	δV_s , km/s
BGCA	5.2	18.4	54	88	242	0.006	465	-0.2
BNG	4.4	18.5	31					
CMB	38.0	239.6	69	77	138	0.007	500	-0.3
TUC	32.3	249.2	98	86	241	0.006	495	-0.1
ALE	9.4	42.0	23	83	86	0.007	510	-0.2
DSS	11.1	39.6	38					
BJT	40.0	116.2	81	81	125	0.004	465	-0.2
YKW3	62.5	245.4	122	85	334	0.005	485	-0.3
SCHQ	54.8	293.2	55	77	181	0.007	470	-0.2

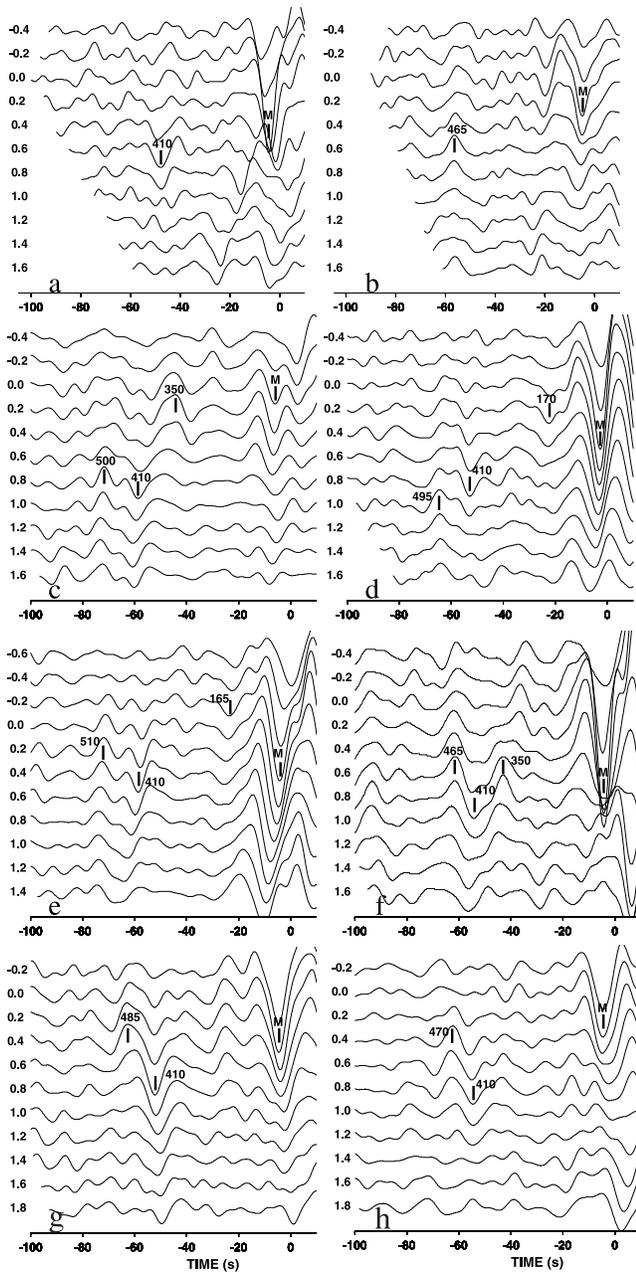


Figure 2. Stacked S receiver functions for seismograph stations BGCA and BNG in back azimuths around (a) 78° and (b) 240° ; stations (c) CMB, (d) TUC, (e) ALE and DSS, (f) BJT, (g) YKW3, and (h) SCHQ. Each trace corresponds to the differential slowness in s° on the left. The arrivals of the Sp phases are marked on the traces with the largest amplitudes; the numbers indicate depths of the discontinuities in km; M is for the Moho. For the details see Table 1.

from the sharp discontinuity. A gradational boundary is suggested by the fact, that signals from the negative discontinuity were never identified in short periods in the PRFs (see Introduction). The thickness of the underlying low velocity layer is usually not well constrained, because signals from its lower boundary are not detected with confidence. Among other reasons, this can be an effect of a reduced sensitivity of the method at larger depths. The model of the MTZ thus

obtained can be described by the two parameters: depth of the middle of the gradational boundary and S velocity reduction in the underlying layer relative to IASP91 (Table 1). The accuracy of depth of the middle of the gradational boundary thus determined is not better than 10–15 km; the variations of this depth between 465 and 510 km (Table 1) may reflect both actual lateral variations and errors of the estimates.

[12] To compare SRFs and PRFs we calculated PRF for station SCHQ in the same azimuth and frequency range as the corresponding SRF (Figure 3b). The time interval of interest between 44 and 70 s in the PRF is dominated by reverberations and no details within the MTZ can be found with confidence. The P660s phase is probably seen, but its theoretical amplitude is at least twice larger than the expected amplitude of the signal from the negative discontinuity. A similar result is obtained for station TUC. Note that the piercing points of the Sp and Ps phases from a discontinuity in the MTZ are at a distance of several hundred kilometers from each other, and they not necessarily sample the same region in the MTZ.

3. Anisotropic Models of the MTZ

[13] Here we try to explain the Sp phase with positive polarity by a mechanism increasing the S velocity over the discontinuity. Among the volumetrically important minerals of the upper MTZ, wadsleyite crystal shows the largest elastic anisotropy, and deformations lead to a preferred crystal orientation [Tommasi *et al.*, 2004]. We investigate a possibility that the Sp phase with positive polarity is generated at the boundary between anisotropic wadsleyite and practically isotropic ringwoodite [Wenk *et al.*, 2005]. Theoretical seismograms are synthesized by using Thomson-Haskell-Crampin algorithm [Kosarev *et al.*, 1979] for elastic parameters of deformed pyrolite in Table 1 of Tommasi *et al.* [2004]. These parameters are obtained for the MTZ pressure at ambient temperature, and therefore the related isotropic wave velocities are higher than the actual velocities. To obtain lower velocities, we modified the model. In the modified model, anisotropy is the same, but the velocities are close to the actual velocities in the MTZ (9.3 km/s and 5.07 km/s for the P and S waves, respectively). The assumed wave slowness is $10 s^\circ$. The boundaries between anisotropic and isotropic pyrolite are placed at 410 km and 520 km depths; elastic parameters of the isotropic media are those from standard Earth models.

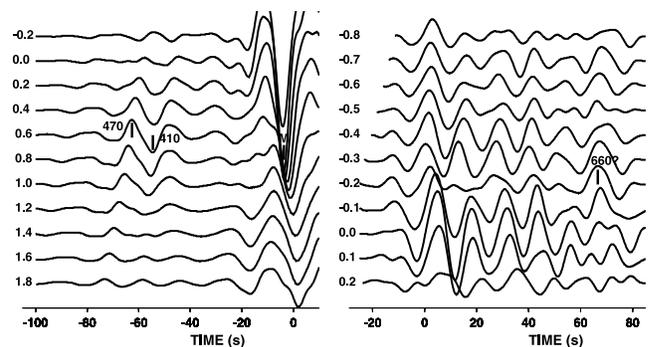


Figure 3. (left) Synthetic SRF and (right) observed PRF for station SCHQ.

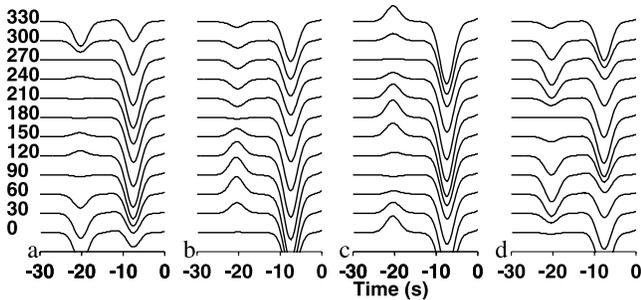


Figure 4. Synthetic P components for deformed wadsleyite: (a) shear in the vertical direction, shear plane is in the azimuth of 90° ; (b) horizontal shear plane, azimuth of shear is 90° ; (c) horizontal foliation plane, azimuth of lineation is 0° ; (d) vertical foliation plane in the azimuth of 0° , lineation is vertical. Origin of the time axis is arbitrary. Numbers attached to the traces are back azimuths in degrees.

[14] The synthetic SRFs (Figure 4) contain the S410p (at around -10 s) and S520p (at around -20 s). The largest amplitudes of the S410p phase are practically those for isotropic medium, and they can be used as a reference. For the vertical shear (Figure 4a), polarity of the S520p phase is opposite to the observed one everywhere, except the two traces, where its amplitude is very small. For the horizontal shear (Figure 4b), the polarity in some azimuths is positive, but the amplitude is less than 30% of that of S410p, much lower than in the actual data (Figure 2). The elastic parameters of Tommasi *et al.* [2004] are for relatively low strain. At larger strains the foliation and shear planes are practically parallel. In the case of horizontal foliation plane (Figure 4c), the polarity in some azimuths is positive, but the amplitude is too small. Moreover, similar amplitudes at stations CMB and TUC in the azimuths differing by roughly 90° contradict the synthetics. Finally, in the case of vertical foliation plane (Figure 4d), polarity of S520p is always negative. We conclude that the synthetics cannot be reconciled with the actual data mainly due to the large difference in amplitudes.

4. Discussion and Conclusions

[15] At several locations we observed seismic phases that, in our interpretation, are converted from a negative discontinuity at a depth between 465 and 510 km. At the negative discontinuity S velocity decreases with depth by about 0.2 km/s. The synthetics for anisotropic wadsleyite do not fit the actual SRF, and thus, anisotropy as an explanation for the negative discontinuity seems unlikely.

[16] Locations of the anomalous regions follow a remarkable regularity. Afar and Iceland are classic hotspots. The piercing points for stations TUC and CMB are in the region of Baja and Guadalupe hotspots, which are often referred as one hotspot. The piercing point for station TUC coincides with the Guadalupe. The piercing point for station CMB is located to the north-east of the Guadalupe, but the first Fresnel zone of the Sp phase in the MTZ is a few hundred kilometers in diameter, and we believe that both stations sample the same anomaly of the MTZ. The piercing point for stations BNG and BGCA is near Cameroon hotspot. This hotspot contains a

number of volcanic centers along the SW-NE line in the ocean and continental margin. The piercing point corresponds to the middle of the line. The region sampled by the data of station BJT is remarkable by rifting, high heat flow, erosion of the mantle lithosphere and low seismic velocities in the upper mantle [Griffin *et al.*, 1998]. These properties are indicative of a mantle upwelling, which might originate in the MTZ. The piercing points for stations YKW3 and SCHQ are near low S velocity anomalies in the MTZ (Figure 1).

[17] We presented practically all available data on the MTZ beneath hotspots, and out of the four sampled hotspots all display the negative discontinuity. The SRF may sample the MTZ beneath a hotspot only if the seismograph station is at the appropriate distance from the hotspot, and the seismic events are in the appropriate back azimuth. For example, station TAM (coordinates 22.8 N, 5.5 E), without any indication of the negative discontinuity in the MTZ, is located in the Hoggar hotspot, but its SRF samples the MTZ ~ 700 km to the west of the station, far from the volcanic centers.

[18] As we already mentioned in the Introduction, if on the average the MTZ is almost dry, a local anomaly of 1 wt% of water may generate a significant S wave velocity anomaly, and reduced S wave velocities are generally indicative of hydration [Smyth and Jacobsen, 2006]. Solubility of water in wadsleyite sharply decreases with temperature [Demouchy *et al.*, 2005], and, if the water content in the MTZ at a low temperature is locally close to saturation, an increase in temperature on the order of 100° may cause dehydration and melting with a large reduction of S velocity. This might explain the connection between the low velocity and hotspots. The melt may reside at a depth around 500 km owing to its neutral buoyancy.

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References

- Demouchy, S., E. Deloule, D. J. Frost, and H. Keppler (2005), Pressure and temperature dependence of water solubility in Fe-free wadsleyite, *Am. Mineral.*, *90*, 1084–1091.
- Duncan, R. A., and M. A. Richards (1991), Hotspots, mantle plumes, flood basalts and true polar wander, *Rev. Geophys.*, *29*, 31–50.
- Farra, V., and L. P. Vinnik (2000), Upper mantle stratification by P and S receiver functions, *Geophys. J. Int.*, *41*, 699–712.
- Griffin, W. L., A. Zhang, S. Y. O'Reilly, and C. G. Ryan (1998), Phanerozoic evolution of the lithosphere beneath the Sino-Korean craton, in *Mantle Dynamics and Plate Interactions in East Asia*, *Geodyn. Ser.*, vol. 27, edited by M. Flower, et al., pp. 107–126, AGU, Washington D. C.
- Hirschmann, M. M., C. Abaud, and A. C. Withers (2005), Storage capacity of H₂O in nominally anhydrous minerals in the upper mantle, *Earth Planet. Sci. Lett.*, *236*, 167–181.
- Kennett, B. L. N., and E. R. Engdahl (1991), Traveltimes for global earthquake location and phase identification, *Geophys. J. Int.*, *105*, 429–465.
- Kosarev, G. L., L. I. Makeyeva, E. F. Savarensky, and E. M. Chesnokov (1979), Influence of anisotropy beneath seismograph station on body waves (in Russian), *Izv. Akad. Nauk, Ser. Fiz.*, *2*, 26–37.
- Ritsema, J., and H. J. van Heist (2000), Seismic imaging of structural heterogeneity in the Earth's mantle: Evidence for large-scale mantle flow, *Sci. Prog.*, *83*, 243–259.
- Smyth, J. R., and S. D. Jacobsen (2006), Nominally anhydrous minerals and Earth's deep water cycle, in *Earth's Deep Water Cycle*, *Geophys. Monogr. Ser.*, edited by S. van der Lee and S. D. Jacobsen, AGU, Washington, D. C., in press.

- Tommasi, A., D. Mainprice, P. Cordier, C. Thoraval, and H. Couvy (2004), Strain-induced seismic anisotropy of wadsleyite polycrystals and flow patterns in the mantle transition zone, *J. Geophys. Res.*, *109*, B12405, doi:10.1029/2004JB003158.
- Vinnik, L. P., V. Farra, and R. Kind (2004), Deep structure of the Afro-Arabian hotspot by S receiver functions, *Geophys. Res. Lett.*, *31*, L11608, doi:10.1029/2004GL019574.
- Vinnik, L. P., G. R. Foulger, and Z. Du (2005), Seismic boundaries in the mantle beneath Iceland: A new constraint on temperature, *Geophys. J. Int.*, *160*, 533–538.
- Wenk, H.-R., G. Ischia, N. Nishiyama, Y. Wang, and T. Uchida (2005), Texture development and deformation mechanism in ringwoodite, *Phys. Earth. Planet. Inter.*, *152*, 191–199.
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