



## 2 North Atlantic geoid high, volcanism and glaciations

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6 [1] Shallow topography, geoid high and intense volcanism  
7 in the northern Mid Atlantic Ridge are interpreted as  
8 enhanced by the loading on the adjacent continents by ice  
9 caps during upper Cenozoic glaciations. The load of ice  
10 packs on the continental lithospheres of North America and  
11 northern Europe generated radial mantle flow at depth. In  
12 our model, these currents, flowing from west and east, faced  
13 each other below the northern Atlantic, joining together and  
14 upwelling. Numerical modeling of this process supports the  
15 development of dynamic topography leading to uplift of the  
16 sea-floor and inducing a regional geoid high. The mantle  
17 rising to shallower levels may have contributed to larger  
18 asthenospheric melting, and to ridge centered excess  
19 magmatism, as observed in the Northern Atlantic.

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### 24 1. Introduction

25 [2] The lithosphere generated by the Mid Atlantic Ridge  
26 (MAR) east of Greenland underlies the youngest (<60 Myr)  
27 and narrowest part of the Atlantic Ocean. This portion of the  
28 northern Atlantic shows three peculiar characters, 1) it is  
29 about 1–3 km shallower than the average mid-oceanic ridge  
30 (Figures 1a and 1b); 2) it displays diffuse positive gravity  
31 (>30 mGal) and geoid (>50 m) anomalies (Figures 1c, 1d,  
32 and 1f); 3) it is the seat of larger than average magmatic  
33 productivity, resulting in the thickest oceanic crust of the  
34 entire MAR, up to about 40 km below Iceland [Kaban *et al.*,  
35 2002]. The thickness of the Cretaceous-Early Cenozoic  
36 (pre-glaciations) oceanic crust in the northern Atlantic is  
37 rather 4–6 km in average [e.g., Shillington *et al.*, 2006].

38 [3] A number of papers attributed these features to the  
39 Iceland mantle plume [Vink, 1984]. However, the deep  
40 hotspot hypothesis has been questioned on various grounds  
41 [e.g., Foulger and Anderson, 2005]: the persistence of  
42 magmatism on the westerly moving ridge and the presence  
43 of a double tail both west and east of Iceland; the absence of  
44 a relevant heat flow positive anomaly, and the possible  
45 presence of a hydrous mantle lowering the melting point  
46 [Bonath, 1990; Asimow and Langmuir, 2003]. There is also  
47 contrasting topologic and tomographic evidence on whether  
48 the source of the plume is in the deep or in the upper mantle  
49 [Foulger *et al.*, 2001; Courtillot *et al.*, 2003; Ritsema and  
50 Allen, 2003; Montelli *et al.*, 2004]. Moreover, the geochem-  
51 ical Icelandic signature is not restricted to Iceland, but  
52 continues both north and south along the Mid Atlantic

Ridge [Taylor *et al.*, 1997]. In this article we test numeri- 53  
cally a model in which the far field superficial loading of the 54  
mantle by the ice caps in North America and northern 55  
Europe can contribute to generate the anomalous features 56  
of the North Atlantic. 57

[4] Jull and McKenzie [1996] and MacLennan *et al.* 58  
[2002] have demonstrated that the removal of ice load over 59  
Iceland triggers volcanism. Here we show that an inverse 60  
correlation can occur for magma production, i.e., the ice 61  
loading on the adjacent continental areas may have contrib- 62  
uted to the uplift of the north Atlantic mantle, to the geoid 63  
anomaly and, possibly, to the higher degree of melting 64  
due to faster adiabatic decompression induced by mantle 65  
upwelling. 66

### 2. Model Description and Results

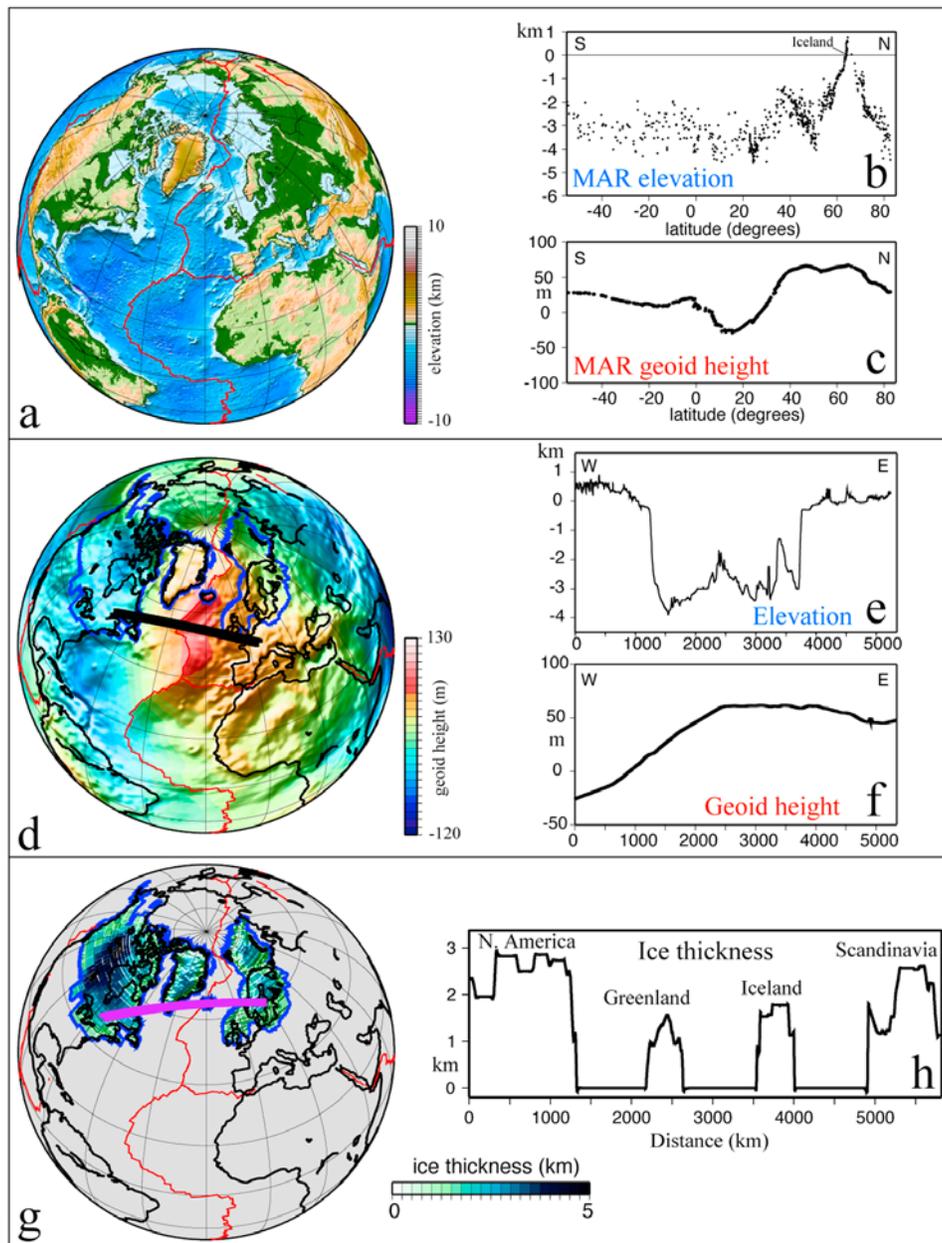
[5] Assuming a viscous Earth (uniform viscous half- 68  
space) and a cylindrical ice-load, it can be shown by 69  
analytical solutions [e.g., Cathles, 1975] that the depth at 70  
which the vertical displacement induced by ice loading/ 71  
unloading is 0.5, 0.2 or 0.1 times the surface value is equal 72  
to  $1.4R_0$ ,  $2.5R_0$  and  $3.3R_0$  (where  $R_0$  is the radius of the 73  
cylinder; i.e., 825 km, 1474 km and 1815 km for the 74  
Fennoscandian ice sheet, characterized by  $R_0 = 550$  km) 75  
respectively. Numerical solutions have also shown that the 76  
ice cycles in the Canadian region induced vertical motions 77  
(either uplift or subsidence) up to more than  $60^\circ$  (more than 78  
6600 km) from the ice center [e.g., Cathles, 1975]. 79

[6] Here we test the combined effects the glacial cycles in 80  
North America and Europe on regional mantle flow. The 81  
aim of our finite element modeling, performed using COM- 82  
SOL 3.5 software (<http://www.comsol.com/>), is to evaluate 83  
the velocity field induced within the upper mantle by 84  
glaciation cycles rather than to reproduce exactly the surface 85  
velocities. This limited objective allowed us to adopt some 86  
major simplifying assumptions, such as the 2D nature of the 87  
model, neglecting the load due to water redistribution 88  
during the ice formation and melting, and using a simplified 89  
ice model. 90

[7] The model adopts a 2D plane strain approximation 91  
and includes lithosphere, upper and lower mantle (Figure 2). 92  
All the layers are described by a compressible linear 93  
viscoelastic (Maxwell) rheology; the assumed elastic con- 94  
stants and viscosities are listed in Table 1. The elastic 95  
structure is consistent with the PREM model [Dziewonski 96  
and Anderson, 1981] and the viscosities are consistent with 97  
values normally used for glacial isostatic rebound modeling 98  
[e.g., Mitrovica and Peltier, 1993; Kaufmann and Lambeck, 99  
2002]. Gravity acceleration and density vary with depth 100  
according to the PREM model. Gravity is applied as a body 101  
force and the ice load as a boundary condition. The ice 102  
thickness varies with time but it is kept laterally constant for 103  
each area. The model is run from 150 Kyr BP to the present. 104

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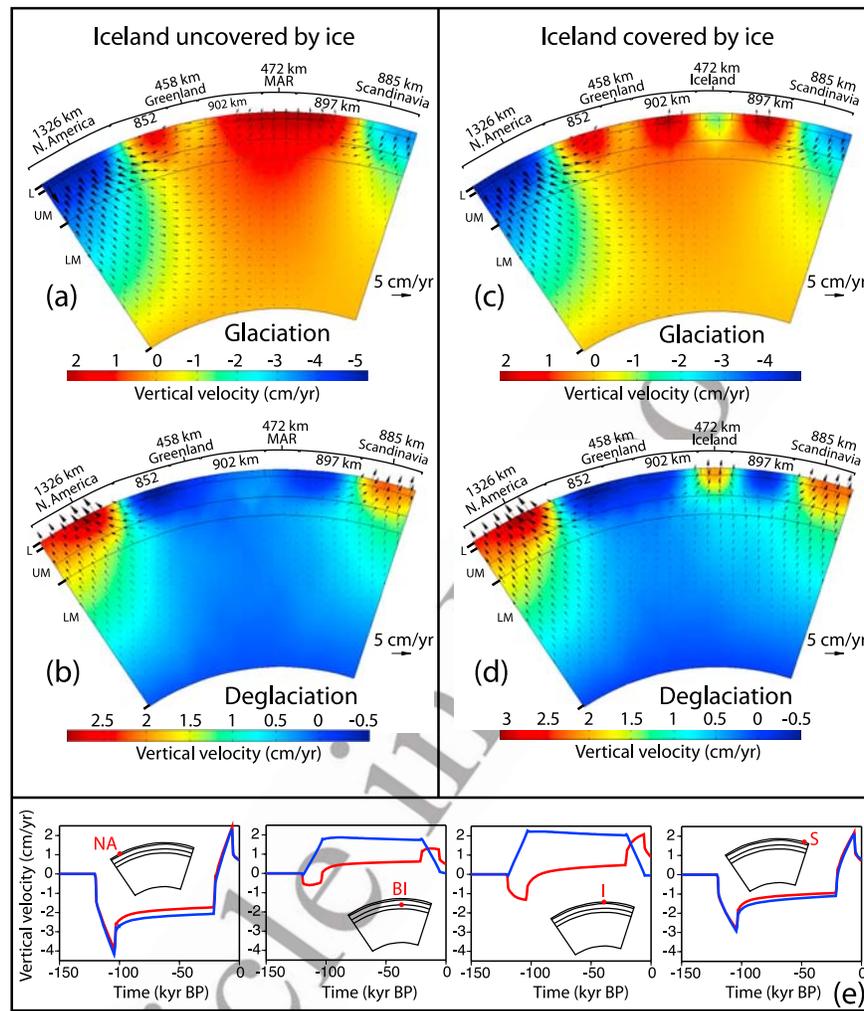
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**Figure 1.** (a) Topography (data after ETOPO1, <http://www.ngdc.noaa.gov/mgg/global/global.html>); (b) elevation of the Mid Atlantic Ridge; the bathymetric distribution along the MAR shows a high in the northern Atlantic which is limited not only to the Iceland area but it extends ca 20° northward and 40° southward; (c) geoid anomaly along the Mid Atlantic Ridge (data after the EGM96 model, <http://cddis.nasa.gov/926/egm96/egm96.html>); (d) geoid height; notice how the northern Atlantic geoid high is located between the North American and Scandinavian ice bodies; (e) topography-bathymetry along the cross-section on the map to the left; (f) geoid height along the same section. The blue curves in panel Figure 1d show the borders of the ice bodies according to ICE-5G. The geoid is shallower along the eastern flank of MAR and the crest of the anomaly is offset to the east of the oceanic ridge. (g) Thickness in map and (h) cross-section in purple of the ice cap at the last glacial maximum (21 Kyr BP; data after the ICE-5G model [Peltier, 2004]). Mid-ocean ridges are shown as red lines. The purple great circle in Figure 1g shows the trace of the modeled profile of Figure 2.

105 The ice thickness is kept at zero between 150 Kyr and  
 106 120 Kyr BP and then it is linearly increased to reach the  
 107 maximum thickness at 105 Kyr BP. It is then kept constant  
 108 until 21 Kyr BP. Between 21 Kyr and 6 Kyr BP the ice  
 109 thickness is linearly decreased to zero, with the exception of  
 110 Greenland, where it is decreased to 750 m. The maximum  
 111 thicknesses is assumed to vary regionally: 2500 m for North

America, 1300 m for Greenland, 2000 m for Scandinavia  
 112 and 2000 m for Iceland (when applied). Such values are  
 113 consistent with the diagram of Figure 1h, showing maxi-  
 114 mum ice thicknesses along the trace of the modeled section  
 115 at 21 Kyr BP according to the ICE-5G model [Peltier,  
 116 2004]. The bottom of the model is fixed normally to the  
 117 boundary and free to slip tangentially. Symmetry conditions  
 118



**Figure 2.** Vertical velocities and velocity fields predicted by the models (a and c) immediately after the formation of the ice caps (105 Kyr BP) and (b and d) soon after their melting (6 Kyr BP). Figures 2a and 2b are referred to a model characterized by the absence of ice during the glacial period in the Iceland region, while panels c and d to a model characterized by Iceland affected by ice. (e) Vertical rates through time at four different locations marked by the red dots (NA, North America; BI, beneath Iceland; I, Iceland; S, Scandinavia). Blue line is with Iceland unaffected by ice, whereas the red line represents the case of Iceland covered by ice.

119 are imposed on the left and right boundaries. This is  
 120 reasonable since the tips of the modeled section are located  
 121 approximately at the center of the American and Scandina-  
 122 vian ice masses. The model surface is left free in areas  
 123 unaffected by ice formation. We use a set of ca. 3800  
 124 triangular elements. Modeling results are shown for the  
 125 time steps of 105 Kyr and 6 Kyr BP, representative of the  
 126 glaciation and deglaciation scenarios respectively. Although  
 127 no constraints are available for past mantle velocity simu-  
 128 lations, we are confident that the patterns and the order of  
 129 magnitude of the calculated velocities are realistic. This  
 130 confidence is justified by the positive match between  
 131 simulated and observed present-day vertical velocities for  
 132 well-constrained areas such as Scandinavia.

133 [8] Two scenarios are modeled. In the first Iceland is  
 134 covered by ice during the glaciation, while in a second  
 135 Iceland is assumed to be ice-free. The first model simulates  
 136 the evolution in the transect of Figure 1, while the second  
 137 simulates a section just north or south of Iceland. Figure 2  
 138 shows the vertical velocities and the velocity field predicted

for the two scenarios. Both scenarios indicate a convergence  
 of velocity vectors towards the Atlantic area during forma-  
 tion of the ice cap, with a prevalence of horizontal directions  
 of motion. Below Iceland and the surrounding Atlantic the  
 velocity vectors turn vertical with a general upwelling (rates  
 of up to 2 cm/yr in the Iceland ice-free scenario). In the  
 Iceland-covered scenario, the upwelling is limited to the  
 Atlantic region with rates of less than 2 cm/yr. Below  
 Iceland the lowermost upper mantle moves upward at slow  
 rates ( $<0.5$  cm), while the shallower upper mantle moves  
 downward, due to the Icelandic ice load. During the same  
 glaciation period, a negative (i.e., downward) velocity field  
 with rates of  $-2/-4$  cm/yr is predicted for North America  
 and Scandinavia.

[9] The velocity field is reversed during deglaciation,  
 with the mantle flowing downward and away from the  
 central Atlantic region and upward below Scandinavia and  
 North America. Figure 2e shows that the development of  
 the velocity field associated to glaciation and its reversal  
 during deglaciation is very fast, due to the elastic compo-

t1.1 **Table 1.** Elastic and Viscous Parameters Used in the Calculations

t1.2 Layer	Poisson's Ratio	Young Modulus (Pa)	Viscosity (Pa s)	Depth Interval (km)
t1.3 Lithosphere	0.27	1.75e11	5e22	0–100
t1.4 Upper Mantle	0.27	1.39e11	1e21	100–670
t1.5 Lower Mantle	0.27	1.27e11	1e22	670–2890

159 nent of rheology. Present-day rates, although with lower  
 160 magnitude, show for the two scenarios velocity patterns  
 161 similar to those of Figures 2b and 2d. This is consistent with  
 162 literature [e.g., Vestøl, 2006; Milne et al., 2001]. Thus the  
 163 dynamic topography attained during the glaciation period  
 164 has not been completely recovered, due to the viscous  
 165 component of the rheology of lithosphere and mantle.  
 166 Although not shown, a sensitivity analysis showed that  
 167 the described patterns of the velocity field are stable also  
 168 when the rheological parameters and ice thickness are  
 169 modified within reasonable bounds.

170 [10] Therefore the models show that the ice load induces  
 171 a upward flow below the Mid Atlantic ridge generating a  
 172 dynamic topography consistent with the geoid high mea-  
 173 sured in the region. The results of the model that assumes  
 174 Iceland free of ice allowed us to predict, at 21 ka BP (i.e.,  
 175 just before the beginning of deglaciation), a geoid anomaly  
 176 of ca. 70 m for the center of the Atlantic ocean (location I in  
 177 Figure 2e). The geoid anomaly was calculated as  $\Delta h =$   
 178  $-\frac{2\pi G}{g} \int \Delta \rho(z) z dz$  [Turcotte and Schubert, 2002], where  
 179  $\Delta h$  is the geoid anomaly,  $g$  is the gravity acceleration,  
 180  $\Delta \rho(z)$  is the anomalous density at depth  $z$  and  $D$  is the  
 181 compensation depth (chosen as the bottom of our model) and  
 182  $G$  is the Newtonian constant ( $6.67 \times 10^{11} \text{ m}^3 \text{ kg}^{-1} \text{ m}^{-2}$ ).  
 183 Although this calculation is to be considered a rough  
 184 estimate, since it includes only the upward motion of  
 185 particles below the MAR and does not include crust  
 186 formation, mantle partial melting and other thermal pro-  
 187 cesses, it is compatible with the present day anomaly of the  
 188 region (ca. 60 m; Figure 1), showing that present-day geoid  
 189 anomaly and high topography of the region are remnants of  
 190 the glaciation. These findings also explain the topographic  
 191 low below Scandinavia and North America, consistent with  
 192 the observed geoid low (the low geoid anomaly of North  
 193 America has been already tentatively explained with the ice  
 194 load by Turcotte and Schubert [2002]).

195 [11] Moreover, mantle upwelling may enhance mantle  
 196 partial melting and explain, at least in part, the anomalously  
 197 intense magmatic activity of the region. Assuming an  
 198 average 7–10% melt of the asthenosphere [e.g., Langmuir  
 199 and Forsyth, 2007] under the northern Mid Atlantic Ridge,  
 200 the cumulative uplift of ca. 2 km of the mantle during the  
 201 glaciations would increase the melting by a few percent  
 202 (depending on water content, initial mantle composition and  
 203 temperature, spreading rate, etc.), producing a larger volume  
 204 of magma delivered to the surface.

### 205 3. Discussion and Conclusions

206 [12] Our modeling has shown, consistently with previous  
 207 studies, that ice loading/unloading can have a regional  
 208 impact on mantle flow velocities. The MAR swollen  
 209 bathymetry (Figure 1) and the geoid regional positive  
 210 anomaly of the northern Atlantic [Lemoine et al., 1998;  
 211 Tapley et al., 2005] are located in an area intermediate

between the ice caps in Northern America and Europe 212  
 during the last glaciation. Moreover, the same area is 213  
 occupied by the largest volcanic province of the northern 214  
 Atlantic. If our model is correct, we speculate a glacio- 215  
 eustatic Milankovitch periodicity in north Atlantic magma 216  
 production. 217

[13] The oldest rocks in Iceland are about 15 Ma old 218  
 [Hardarson et al., 1997]. The same Authors noted chemical 219  
 variations of basalts, generated by a variably depleted 220  
 mantle. Iceland possibly emerged at that time or later, and 221  
 it experienced ice loading as well. The time of the onset of 222  
 glaciations in the northern hemisphere is still debated. It has 223  
 been shown how the onset of glaciations in the northern 224  
 hemisphere is older (Eocene-Oligocene) than previously 225  
 estimated [Eldrett et al., 2007]. Recent deep sea drilling 226  
 provided evidence for a middle Eocene initiation of the 227  
 icehouse of the Arctic area [Moran et al., 2006]. High 228  
 magma productivity has been documented in Iceland 13– 229  
 11 Myr, and 8–7 Myr intervals together with periodicity in 230  
 magma composition [e.g., Kitagawa et al., 2008]. Gee et al. 231  
 [1998] detected a close relationship between the geochem- 232  
 istry of lavas and glacioisostasy. They found that eruption 233  
 of primitive lavas with depleted chemical and isotopic 234  
 characteristics coincides with a period of glacioisostatic 235  
 instability at the end of the last glaciation (13–9 Kyr). 236

[14] Sigvaldason [2002] described a Holocene rhyolitic 237  
 eruption triggered by the melting of the ice cap in central- 238  
 eastern Iceland, hinting at a relation between magmatic 239  
 emplacement and vertical loading. 240

[15] Therefore, loading and unloading of the ice cap 241  
 [Watts, 2001] appears to be a factor controlling locally or 242  
 even regionally the production of mantle melts. Although 243  
 we modeled a single ice cycle, the productivity of magma 244  
 over geological periods is expected to be influenced by the 245  
 superposition of several ice cycles on the process of oceanic 246  
 spreading. The remote loading of ice can determine an 247  
 upwelling of the mantle elsewhere, generating larger vol- 248  
 umes of melt due to mantle adiabatic decompression below 249  
 the ridge. Vice versa, the ice load in a volcanic area (e.g., 250  
 along the MAR in Iceland) can locally buffer eruption, 251  
 tuning the frequency of magmatic delivery, and generating a 252  
 lower degree of melting and a longer residence time of 253  
 melts in the mantle. These factors, together with the variable 254  
 source depth of the melts, could cause significant variations 255  
 of the lava's geochemistry. Therefore, in Iceland, the fol- 256  
 lowing two complementary processes could interfere, over- 257  
 lap, and buffer each other: deglaciation-induced magmatism 258  
 (a in-situ mechanism associated with stress release related to 259  
 ice unloading) and glaciation-induced magma production 260  
 (a far-field effect, as shown by our model). In the remaining 261  
 areas of the MAR, not directly covered by ice, a different 262  
 time correlation between magma production and eruption is 263  
 expected. 264

[16] Our model predicts a relatively low intensity of 265  
 magmatism along the northern segment of the MAR during 266  
 the present interglacial period. We note that the North 267  
 Atlantic geoid height is presently decreasing, while it is 268  
 increasing on the adjacent continental areas, as shown by 269  
 the Grace project data [e.g., Tapley et al., 2004]. The 270  
 decrease of the geoid has been related to the melting of 271  
 ice in Greenland [Ramillien et al., 2006], but it could be 272  
 related also to the decreasing upwelling beneath the north- 273

ern MAR due to the absence of ice caps on the continents. Conversely, the continental areas show an increase of the geoid because the mantle is rising, recovering the subsidence previously generated by the ice loading. However, when the mantle rises and melts beneath a ridge [McKenzie, 1984], it becomes lighter [Oxburgh and Parmentier, 1977]. Therefore the process is possibly not entirely reversible since the uplifted and depleted mantle cannot be re-pulled down at its original position, by the down-flow motion induced by deglaciation, because of the permanent increase in buoyancy characterizing the mantle after melting.

[17] During the time frame considered (e.g., say the last 20–30 Ma) we may expect about 180–250 oscillations associated to the eccentricity of the Earth's orbit, or more than twice oscillations in case of obliquity related cycles. The model presented rather shows the effects of only one single cycle of loading and unloading. Assuming an irreversible component on each cycle, the present geoid high would represent the sum of the all episodes, a sort of vibration generating hysteresis in the uplift of the mantle.

[18] In summary, we suggest that the ice caps on the continents of the northern hemisphere generated a flow in the underlying mantle that converges in the northern Atlantic from west and east, upwelling along the northern MAR. The eastward offset of the geoid high relative to the MAR could be due to a larger ice load on the northern American continent, although we cannot neglect a contribution from the relative eastward mantle flow implicit in the notion of the net rotation of the lithosphere [Gripp and Gordon, 2002], able to generate an asymmetry of ocean ridges worldwide [Doglioni et al., 2003]. This model implies that the over production of magmatism in the northern Atlantic could be sourced by the shallower location of the asthenosphere, being the upraise of the asthenosphere pumped from the deep mantle flow.

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